

Nonequilibrium field theory and its applications

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Outline

- Motivation
- The closed time path formalism
- Green's functions
- Applications

Motivation

- Imaginary-time formalism:
 - t traded in for T
 - Static, equilibrium properties
- Study of dynamical systems:
 - Evolution of early universe
 - Chiral symmetry breaking

Closed time path formalism

$$\rho(t) = \sum_n p_n |\psi_n(t)\rangle \langle \psi_n(t)|$$

Quantum Liouville equation

$$i \frac{\partial \rho(t)}{\partial t} = [H, \rho(t)]$$

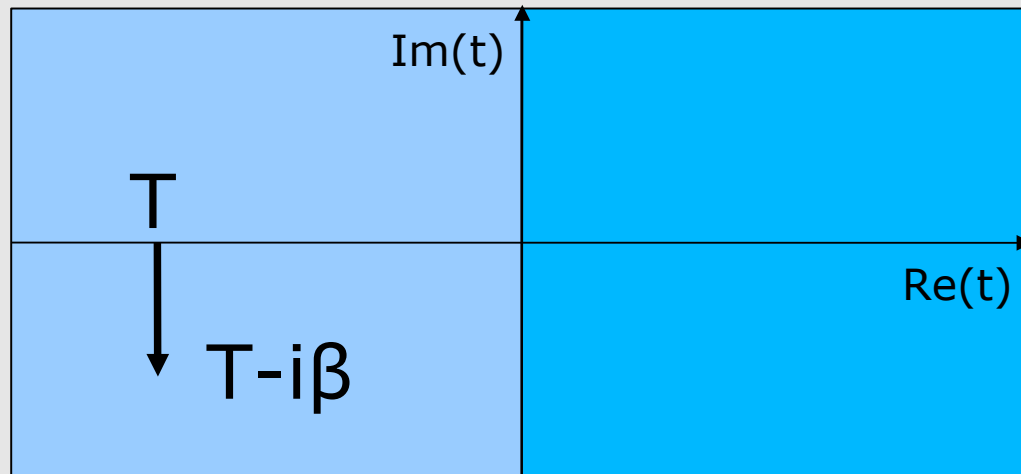
General solution: $\rho(t) = U(t, 0) \rho(0) U(0, t)$

Closed time path formalism

$$H(t) = H_i, \Re(t) \leq 0 \quad \text{Equilibrium state, } \beta$$

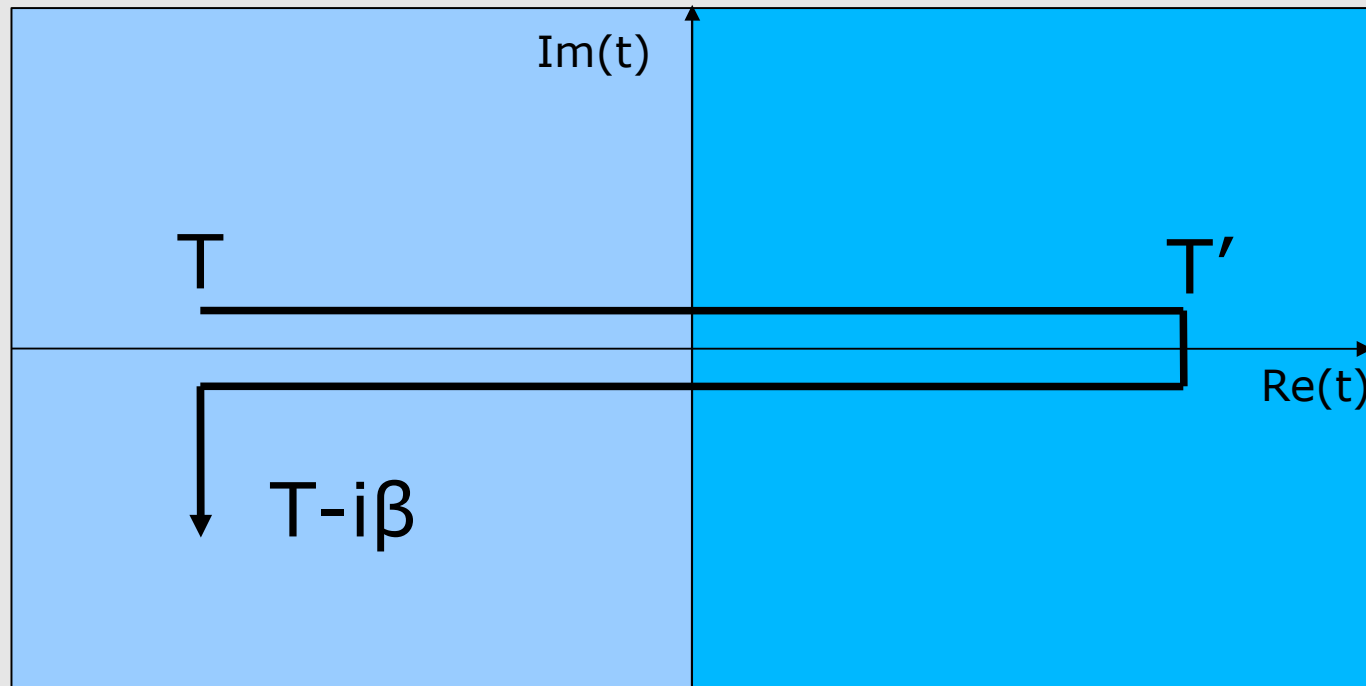
$$H(t) = H(t), \Re(t) > 0$$

$$\rho(0) = \frac{e^{-\beta H_i}}{\text{Tr} e^{-\beta H_i}} = \frac{U(T - i\beta, T)}{\text{Tr} U(T - i\beta, T)}$$



Closed time path formalism

$$\begin{aligned} \langle A \rangle(t) &= \text{Tr} \rho(t) A \\ &= \frac{\text{Tr} U(T - i\beta, T) U(T, T') U(T', t) A U(t, T)}{\text{Tr} U(T - i\beta, T) U(T, T') U(T', T)} \end{aligned}$$



Closed time path formalism

Define generating functional as:

$$Z[J_c] = \text{Tr} U_{J_c}(T - i\beta, T) U_{J_c}(T, T') U_{J_c}(T', T)$$

Or equivalently, setting T, T' to infinity:

$$Z[J_c] = \int D\varphi e^{i \int_c dt \int d^3x (L + J_c \varphi)}$$

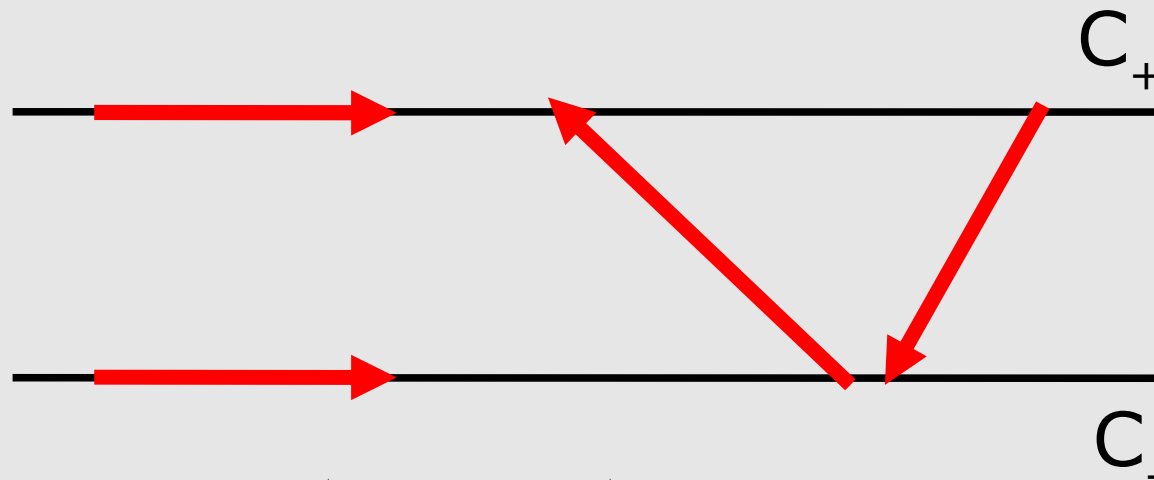
T_c ordered Green's functions

$$G_c(t - t')$$

via functional derivative

Green's functions

Four propagator structures are possible:



$$iG_{++}(t-t') = \langle T(\varphi(t)\varphi(t')) \rangle$$

$$iG_{+-}(t-t') = \langle \varphi(t')\varphi(t) \rangle$$

$$iG_{-+}(t-t') = \langle \varphi(t)\varphi(t') \rangle$$

$$iG_{--}(t-t') = \langle T_{inv}(\varphi(t)\varphi(t')) \rangle \quad \text{anti time ordered!}$$

Green's functions

Double fields and sources:

$$\varphi \rightarrow (\varphi_+, \varphi_-)$$

$$J \rightarrow (J_+, J_-)$$

The action becomes:

$$S = \int d^4x \left[L(\varphi_+, J_+) - L(\varphi_-, J_-) \right]$$

Calculate G in the "standard" fashion:

$$iG_{ab}(x, y) = (-i)^2 \frac{1}{Z} \frac{\delta^2 Z[J]}{\delta J_a(x) \delta J_b(y)} \Bigg|_{J=0}$$

Green's functions

Physical Green's functions in momentum space:

$$G_R(x, x') = \theta(t - t') \langle [\varphi(x), \varphi(x')] \rangle = \frac{1}{p^2 - m^2 + i\epsilon p^0}$$

$$G_A(x, x') = \theta(t' - t) \langle [\varphi(x), \varphi(x')] \rangle = \frac{1}{p^2 - m^2 - i\epsilon p^0}$$

$$G_C(x, x') = \langle [\varphi(x), \varphi(x')]_+ \rangle = -2i\pi (1 + 2n_B(\omega_p)) \delta(p^2 - m^2)$$

G_R and G_A are T-independent at lowest order!

Green's functions

And in Nonequilibrium?

Simply replace

$$n_B \rightarrow f_B$$

in the bare propagators

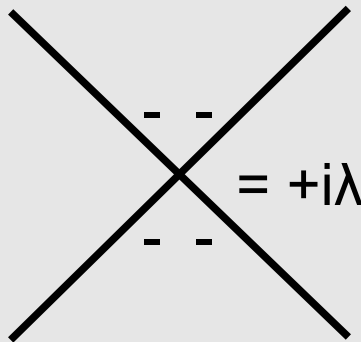
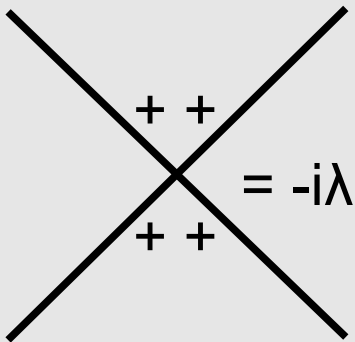
Propagators (one loop)

Consider φ^4 theory:

$$L(\varphi) = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - \frac{m^2}{2} \varphi^2 - \frac{\lambda}{4!} \varphi^4$$

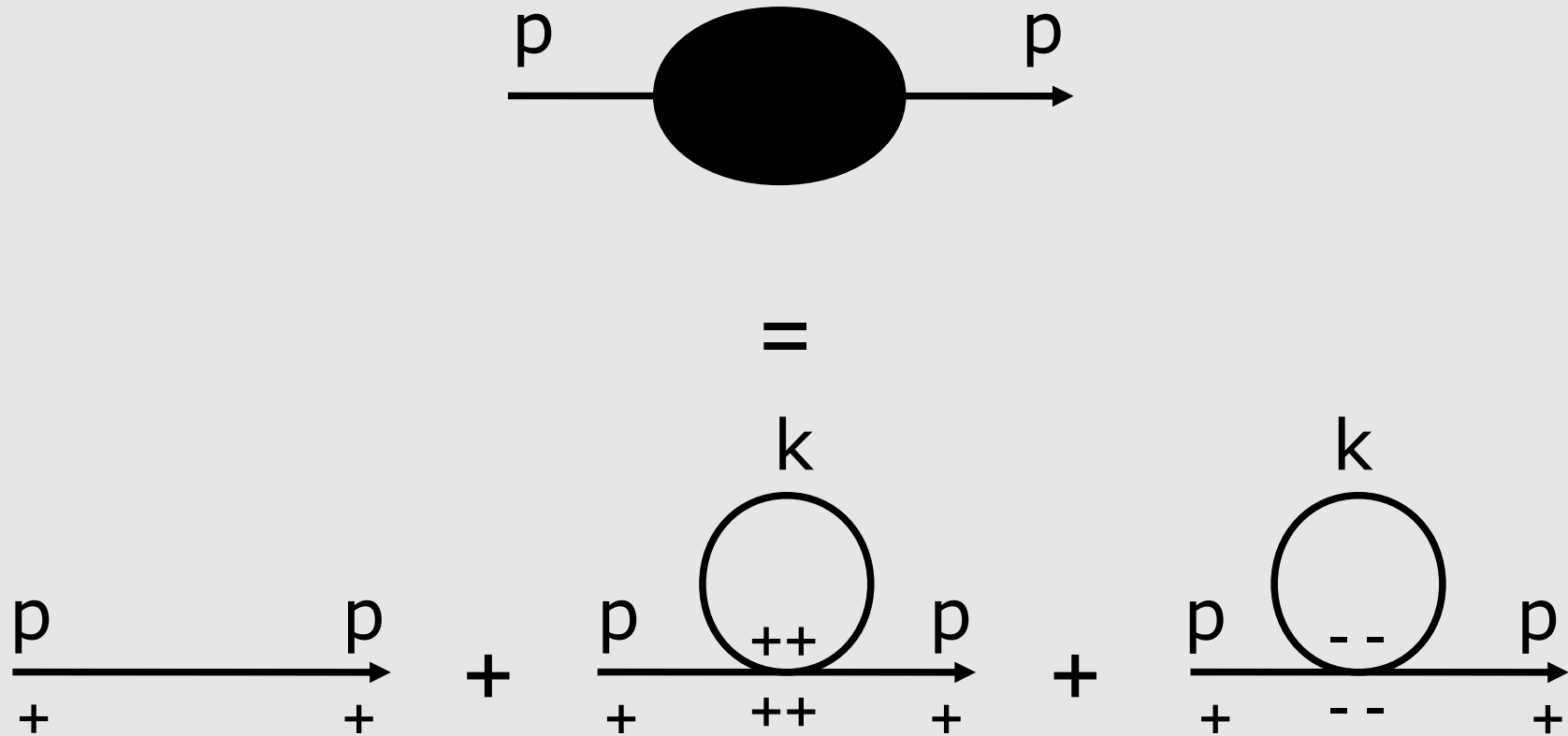
$$S[\varphi_+, \varphi_-] = \int d^4x [L(\varphi_+) - L(\varphi_-)]$$

Two interaction vertices:



Propagators (one loop)

Calculate one loop correction:



Applications

- Domain growth
- Dissipation and fluctuation

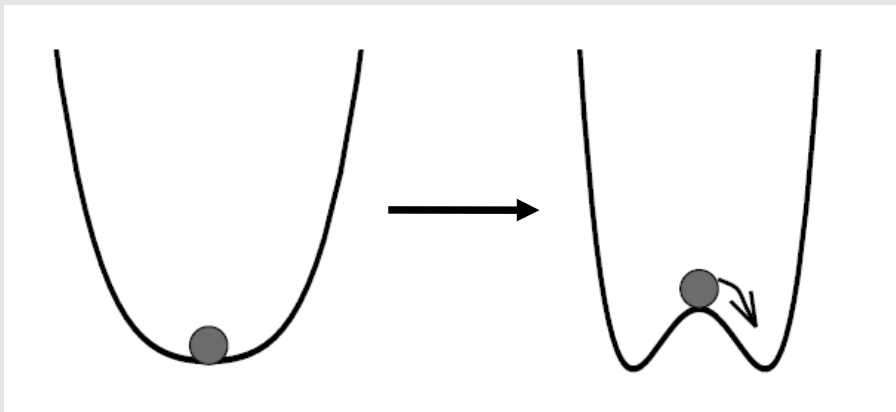
Domain growth

Consider nonequilibrium σ -model

$$L = \frac{1}{2} \partial_\mu \varphi_a \partial^\mu \varphi_a - \frac{1}{2} m_i^2 \varphi_a \varphi_a + \frac{1}{2} \delta m^2(t) \varphi_a \varphi_a - \frac{\lambda}{4} (\varphi_a \varphi_a)^2$$

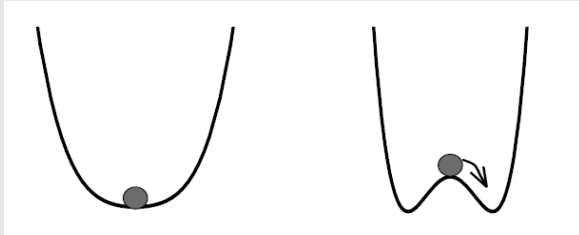
With time dependent mass squared

$$\delta m^2(t) = \theta(t) (m_i^2 + m_f^2)$$



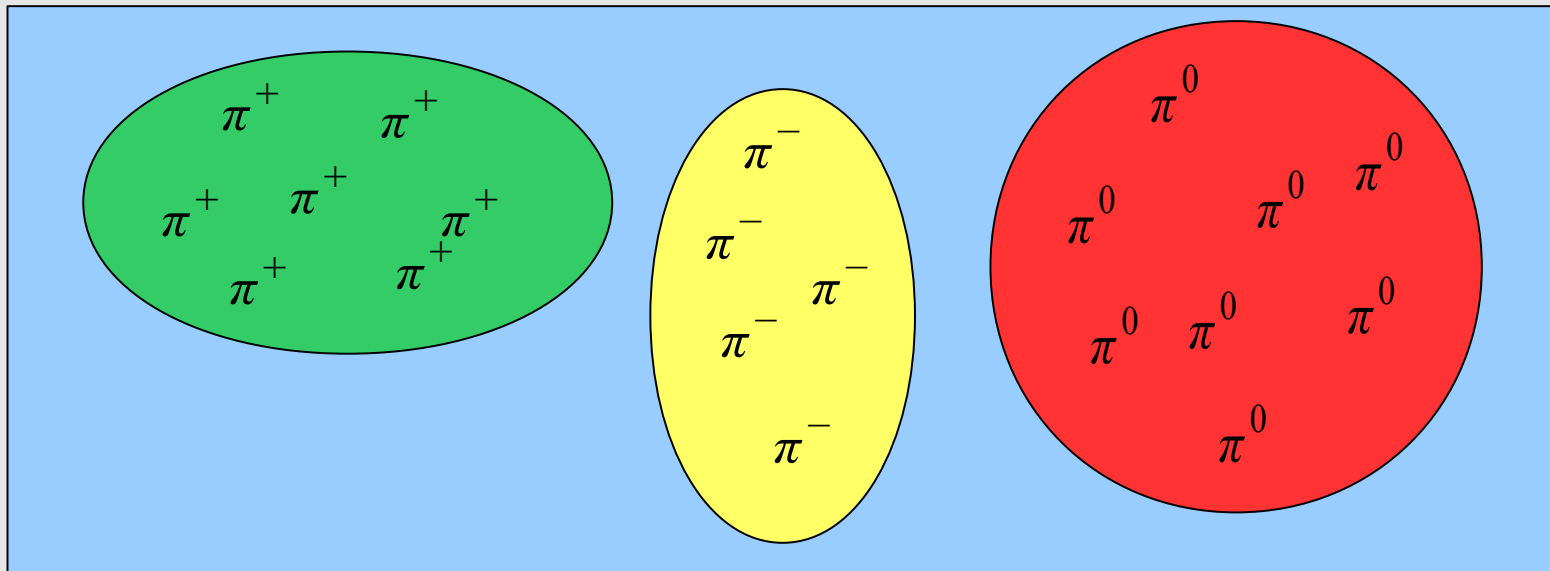
quench at $t=0$

Domain growth



quench at $t=0$

Instability creates long-range correlation,
Disoriented Chiral Condensate



Domain growth

Aim: calculate time evolution of σ -/ π -correlated regions:

$$\langle \sigma(t, \vec{x}) \sigma(t, 0) \rangle$$

For NJL model, 1/N-expansion:

$$G_C(t, |\vec{x}|) = C(T) \left(\frac{|\vec{x}|}{t} \right)^\alpha e^{Mt} e^{\frac{-|\vec{x}|^2}{L^2(t)}}$$

With typical domain size

$$L(t) = \sqrt{\frac{b}{2M}} t$$

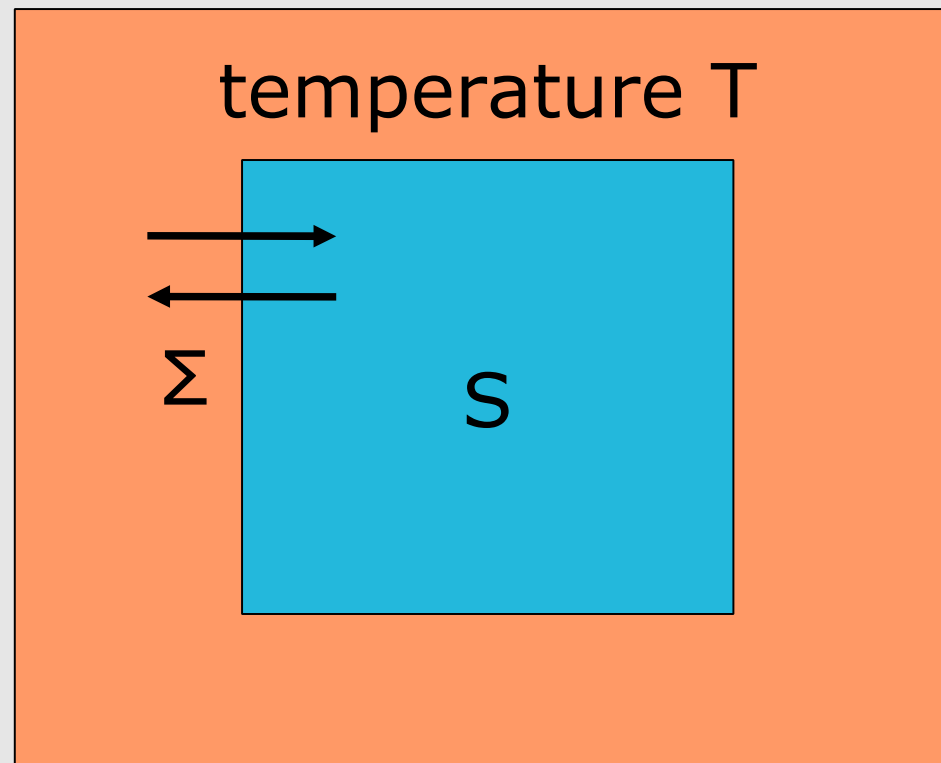
Domain growth

$$L(t) = \sqrt{\frac{b}{2M}} t$$

- Domain size grows linearly with time
- Plasma produces large number of coherent pions
- Indefinite growth due to lack of collisions

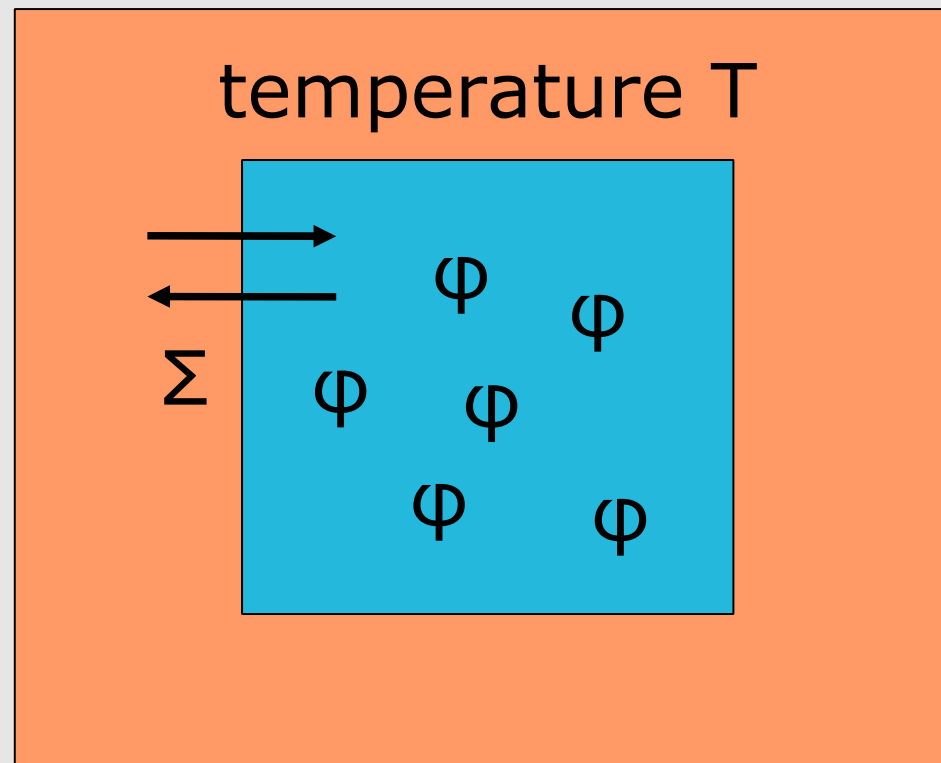
Fluctuation and dissipation

Quantum mechanical system S
interacting with heat bath at temperature T



Fluctuation and dissipation

$$S = \frac{1}{2} \left[\phi_c \left(-\partial_\mu \partial^\mu - m^2 \right) \phi_c - \phi_c \Sigma^c \phi_c \right]$$



Fluctuation and dissipation

Equations of motion for the GFs yield

$$\left(-\partial_{\mu}\partial^{\mu}-m^2\right)G_{<}-\Sigma_R G_{<}-\Sigma_{<}G_A=0$$

Kadanoff-Baym equation

Determines complete and causal non-equilibrium evolution for the two-point functions

Fluctuation and dissipation

Linear combination of $\Sigma_{<}$ and $\Sigma_{>}$ gives

$s(x_1, x_2)$	mass shift
$a(x_1, x_2)$	dissipation term
$I(x_1, x_2)$	fluctuation term

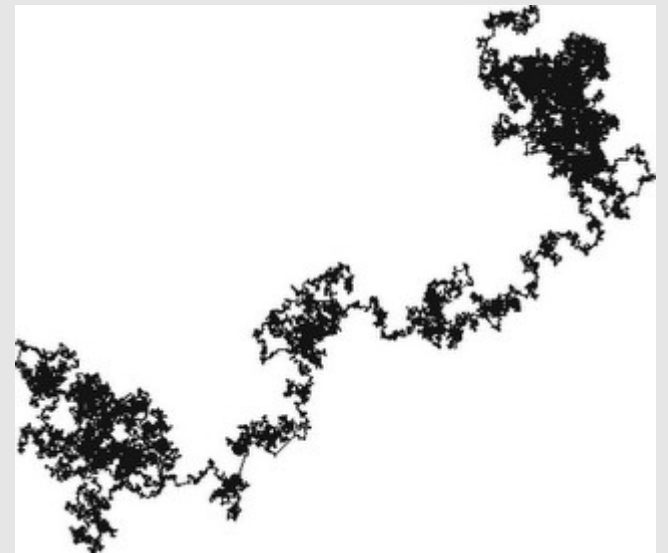
Fluctuation and dissipation

Fourier transformed stochastic EoM for ϕ_ξ

$$\ddot{\phi}_\xi(k, t) + (k^2 + m^2 + s) \phi_\xi(k, t) + 2 \int_{-\infty}^t dt' \Gamma(t - t') \phi_\xi(t') = \xi(k, t)$$

Cf. Brownian motion

$$M \ddot{x} + M\omega^2 x + 2 \int dt' \Gamma(t - t') \dot{x}(t') = \xi(t)$$



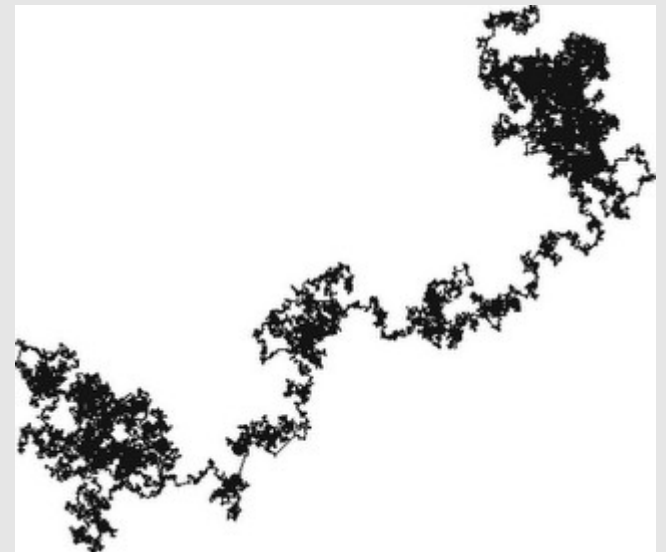
Fluctuation and dissipation

Generalized (microscopic)

Fluctuation-Dissipation-Theorem (high T limit)

$$I(k) = \frac{T}{k_0} 2\text{Im} a(k)$$

- Noise term „heats“ system
- Dissipative term counteracts
- Modes thermally populated
- System will thermalize at T



References

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C. Greiner, S. Leupold: *„Stochastic interpretation of Kadanoff-Baym Equations and their relation to Langevin processes“* (1998)

Thank you